

# Can you hear the shape of Schrödinger's cat?

Colloquium at the Department of Physics of the University of  
Illinois, December 4, 2008

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Department of Physics  
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- ▶ Topological quantum computing?
- ▶ Conclusions



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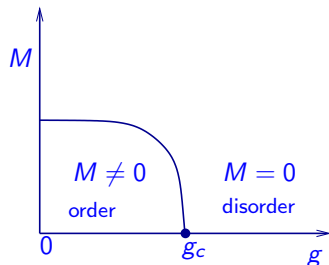
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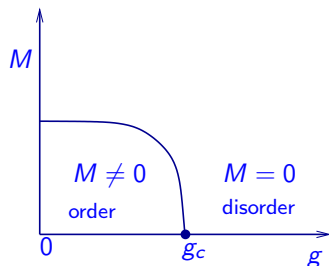


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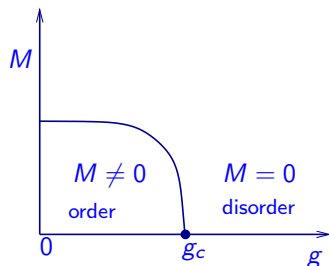
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- ▶ Order parameter:  $M = \langle \sigma_z \rangle$   
 $\Rightarrow$  *spontaneous symmetry breaking*
- ▶ Quantum Critical Point  $g_c$ :  
*scale invariance*, correlation length:  $\xi \sim |g - g_c|^{-\nu}$ ,  
energy gap:  $m \sim |g - g_c|^{z\nu}$   
 $z = 1$  for the Ising model

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- ▶ Natural candidate: Quantum Entanglement, measured by the von Neumann Entanglement Entropy (to be defined shortly...)

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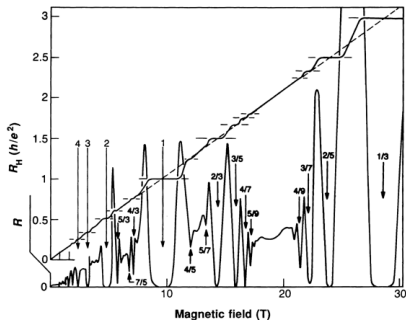
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- ▶ Effective field theory description: Topological Field Theory, e.g., Chern-Simons gauge theory, discrete gauge theory.

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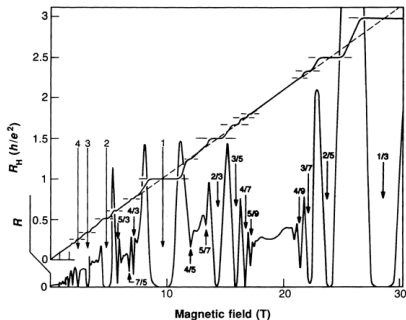
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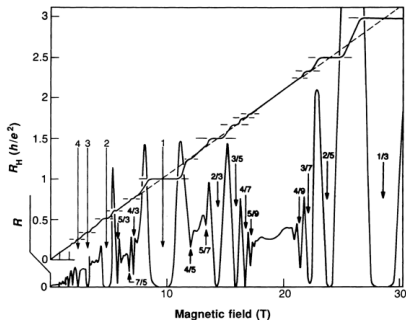
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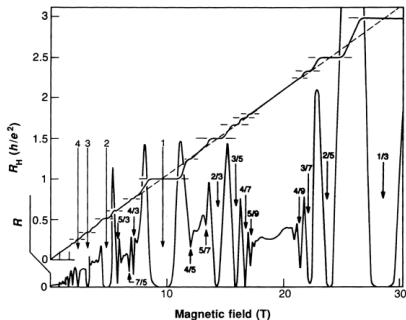


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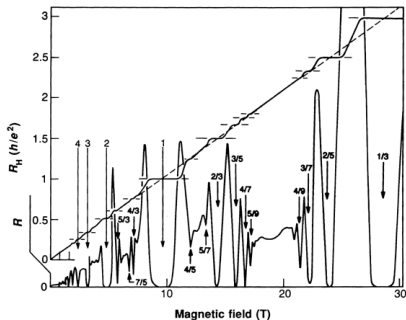
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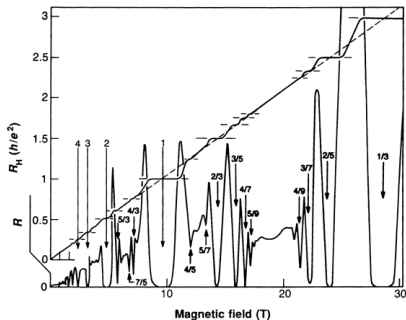
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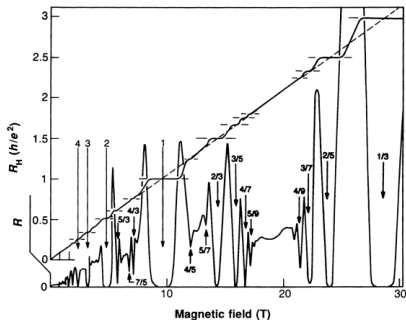
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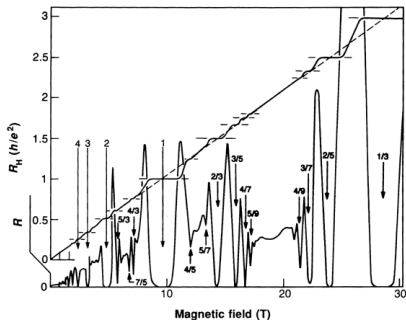
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- ▶ Time-Reversal Breaking Superconductors:  $\text{Sr}_2\text{RuO}_4$  is a  $p_x + ip_y$  superconductor (strong evidence, not controversial)

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  - ▶ Laughlin vortices with charge  $e/m$  and abelian fractional statistics  $\pi/m$

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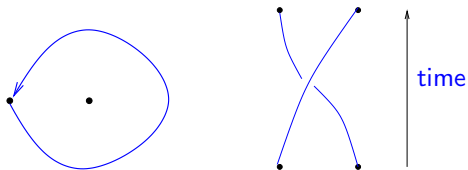
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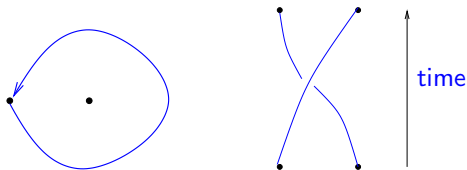
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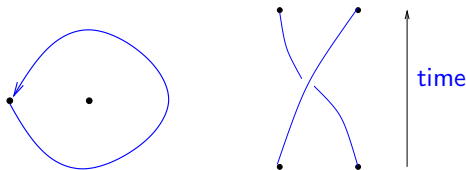
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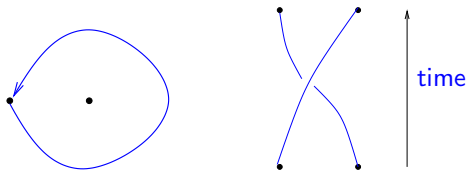
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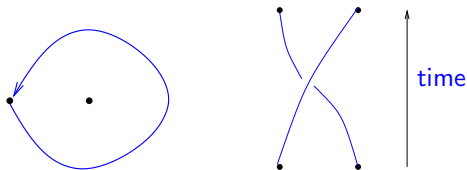


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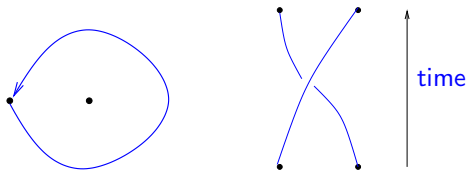
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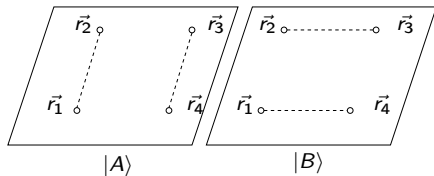
**Anyons** with Abelian (braid) fractional statistics!

# Non-Abelian Braiding Statistics of Quasiholes

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Multivalued wave functions not uniquely determined by the particle coordinates:  
Two linearly independent states at fixed positions of the particles

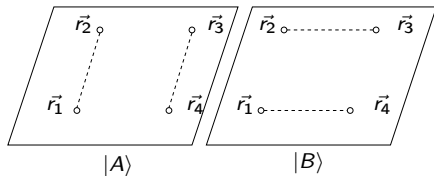


Upon braiding particles 1 and 2  
 $|A\rangle \rightarrow a|A\rangle + b|B\rangle$

# Non-Abelian Braiding Statistics of Quasiholes



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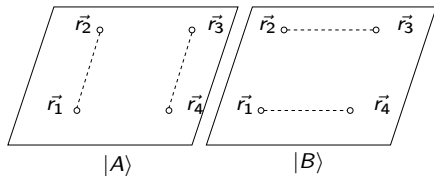
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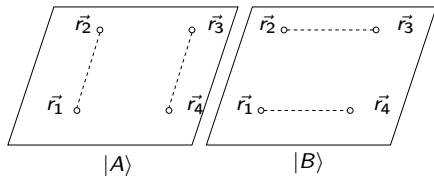
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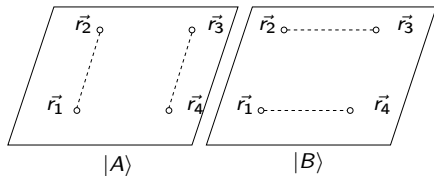
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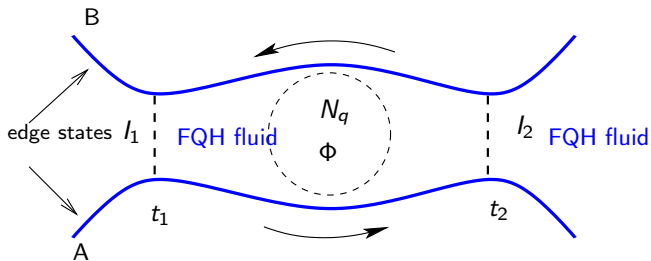
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$$\text{Braiding matrix} \Rightarrow \frac{e^{i\frac{\pi}{4}}}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix}$$

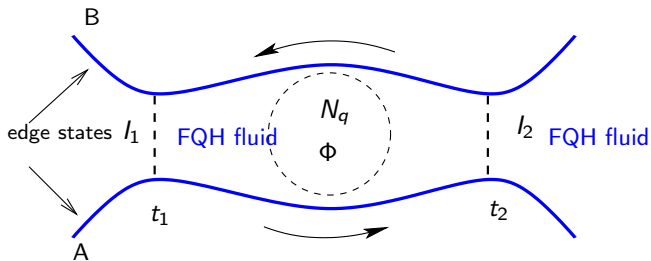
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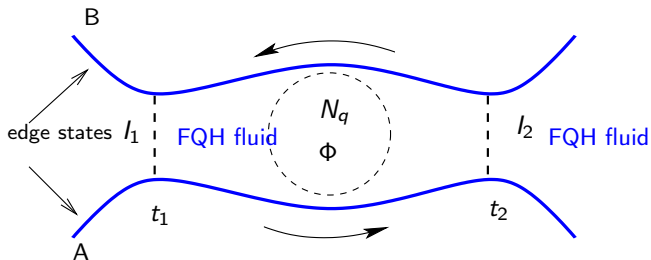
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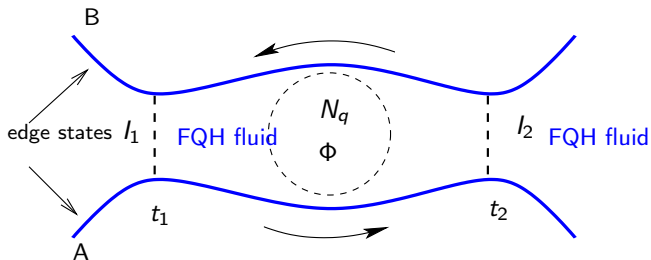
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- ▶ If, on the other hand, we vary  $N_q$ , we can probe the statistics.

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- ▶ Basis of current proposals (Das Sarma, Freedman and Nayak) to construct a “topological qubit”

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- ▶ square lattice: quantum critical point; triangular lattice: Topological  $\mathbb{Z}_2$  deconfined phase (Moessner and Sondhi, 1998)

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- ▶ Matching the correlation functions of the Quantum Dimer and Lifshitz models, one finds  $\kappa = \frac{1}{4\pi}$ .



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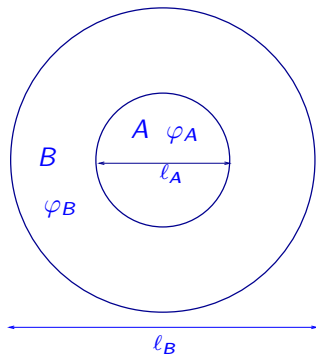
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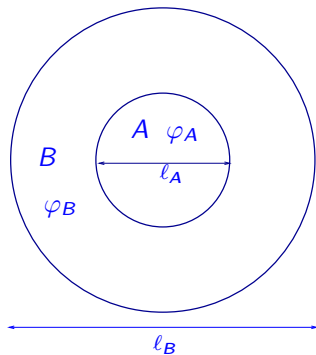
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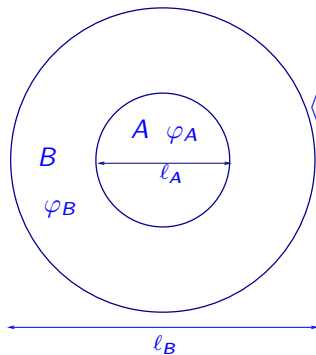
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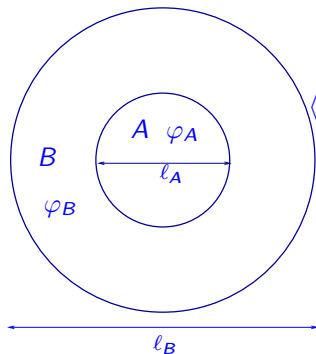


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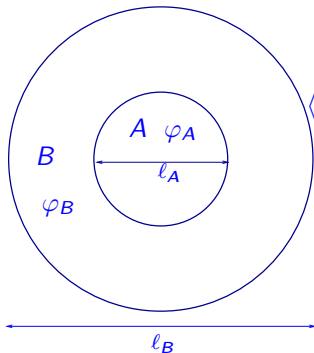
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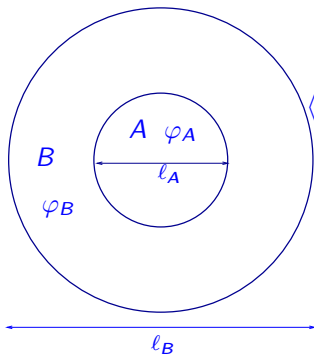
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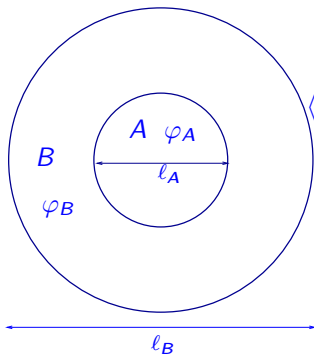
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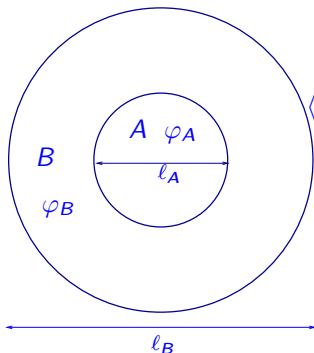
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- ▶ For Fermi liquid in  $d$  dimensions (gapless but not critical),  $S \sim \ell^{d-1} \ln(\ell/a)$  (Wolff, Klich)

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# Entanglement Entropy of Scale Invariant Wave Functions: Can you hear the shape of Schrödinger's cat?

with Joel Moore

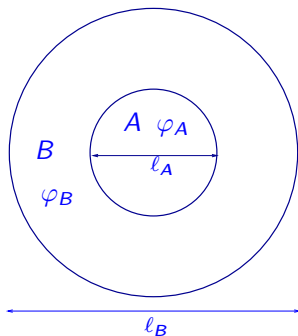
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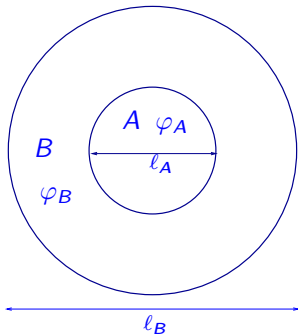
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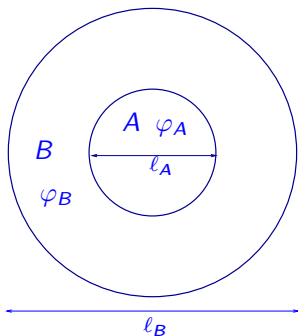
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- ▶  $Z_D^A = \|\Psi\|_A^2$ ,  $Z_D^B = \|\Psi\|_B^2$  and  $Z_{AUB} = \|\Psi\|_{AUB}^2$  with Dirichlet (fixed) boundary conditions

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- ▶ If the  $\ln \ell$  term cancels  $\Rightarrow$  the  $O(1)$  term is universal and determined by the CFT

$$S = \mu \ell + \gamma_{QCP}$$

(There are also geometry-dependent scale invariant terms)  
Hsu, Mulligan, Fradkin and Kim (2008)

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- ▶ In the case of the FQH states this was computed directly from Chern-Simons gauge theory
- ▶ It may be possible to determine the structure of the topological field theory by means of entanglement entropy measurements